

The Basic Physics
of
ELECTROMAGNETICS
WITHOUT
ABSTRACT MATHEMATICS

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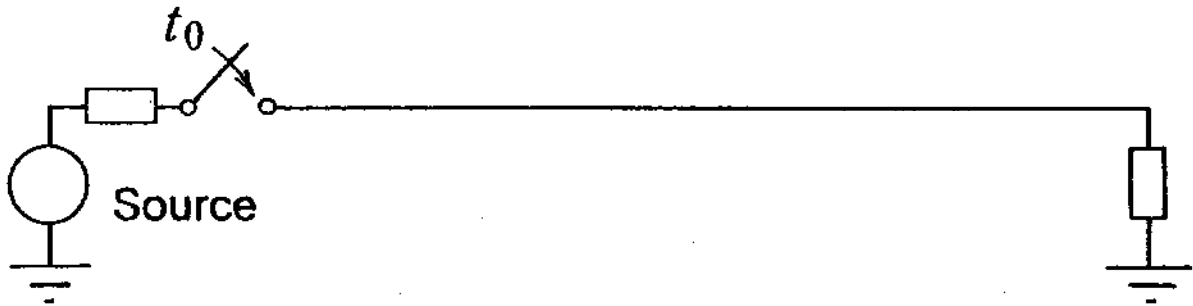
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**– THE CAUSES OF ELECTROMAGNETIC
FIELDS ARE ELECTRIC CURRENTS**

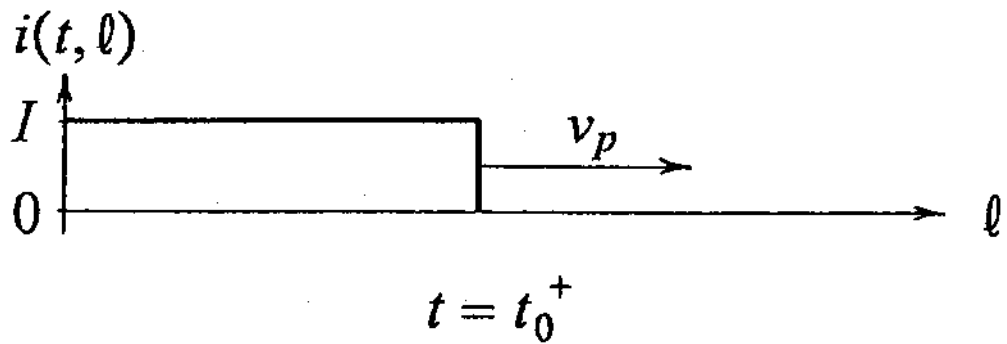
**– A TIME-VARYING ELECTRIC CURRENT
DIFFERS IN VALUE FROM ONE POINT
TO ANOTHER ALONG ITS PATH, AS A
RESULT OF PROPAGATION**

**– THEREFORE, EVERY POINT ALONG A
TIME-VARYING ELECTRIC CURRENT'S
PATH IS A FUNDAMENTAL SOURCE OF
ELECTROMAGNETIC FIELDS**

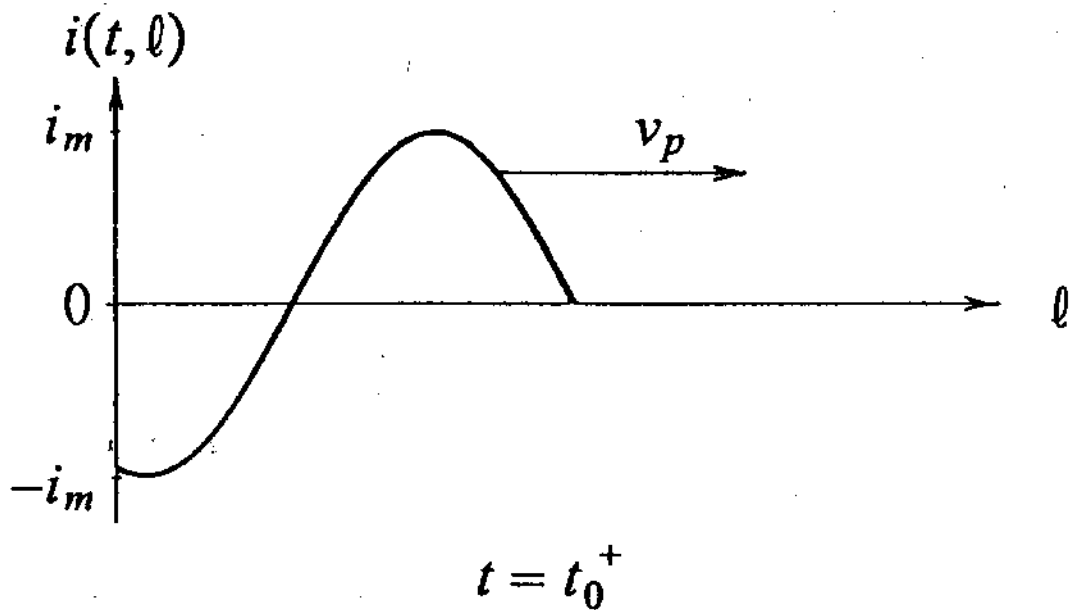
Point-to-Point Circuit-Current Variations



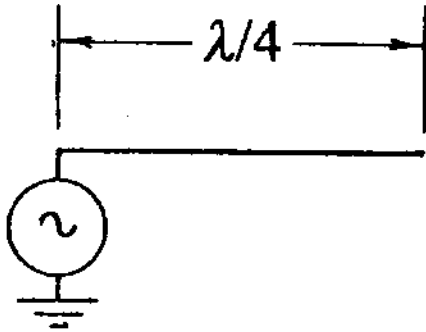
DC Source



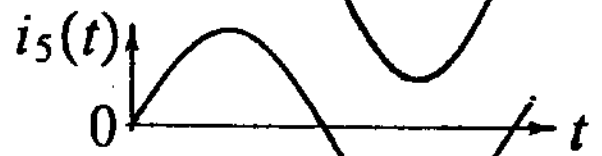
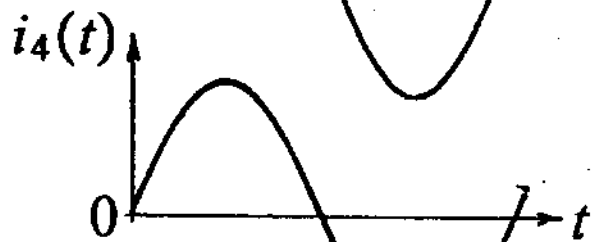
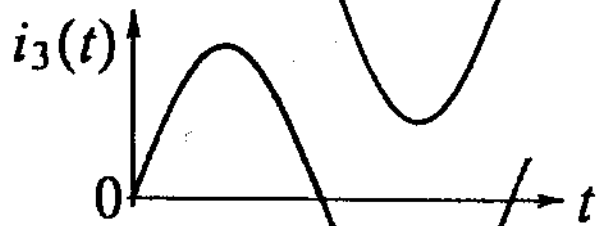
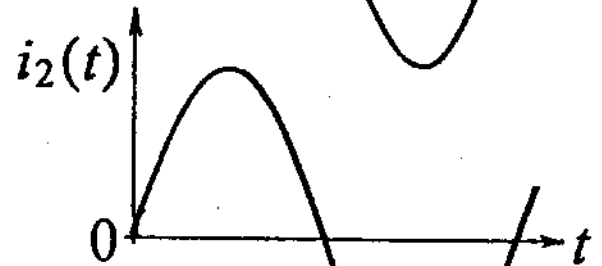
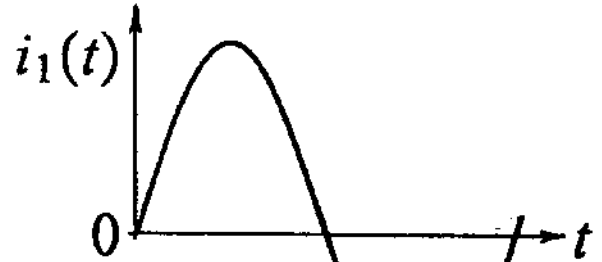
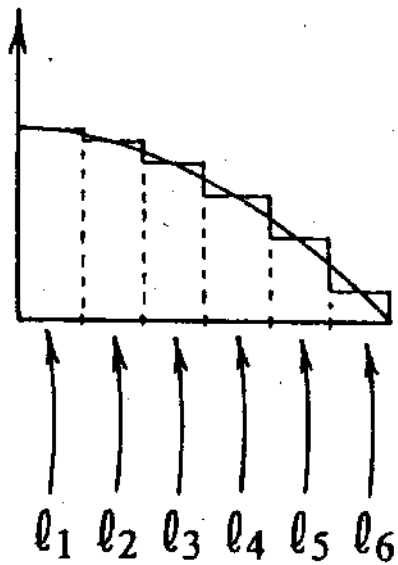
AC Source



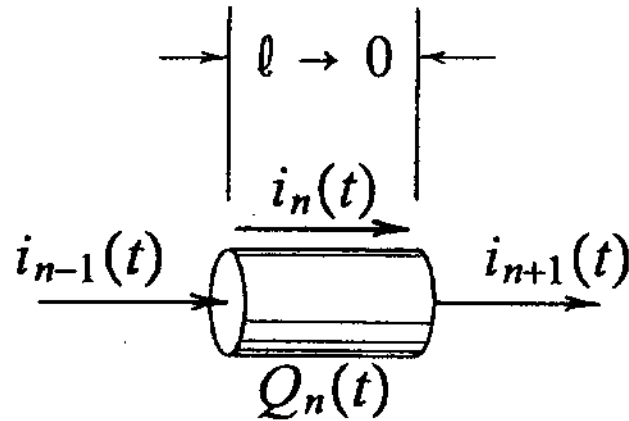
Point Current Variations: $\lambda/4$ Monopole



$|i(t)|$



THE CHARGE ELEMENT



Point Current:

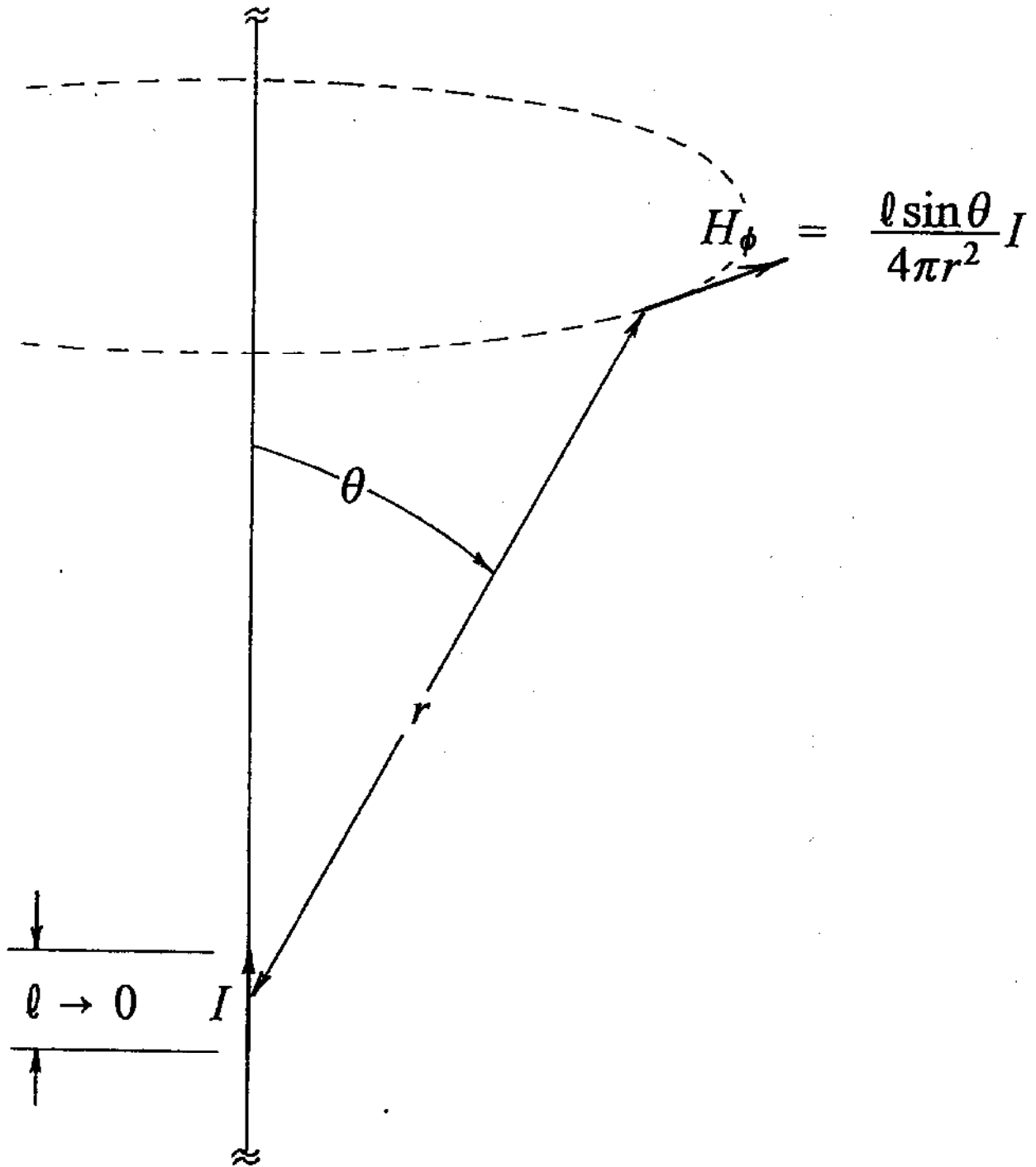
$$\begin{aligned} \ell i_n(t) &= \ell \frac{dq_n(t)}{dt} = \frac{d}{dt} [\ell q_n(t)] \\ &= q_n(t) \frac{d\ell}{dt} = q_n(t) \bar{v} \end{aligned}$$

Point Charge:

$$Q_n(t) = \int_0^t [i_{n-1}(t) - i_{n+1}(t)] dt$$

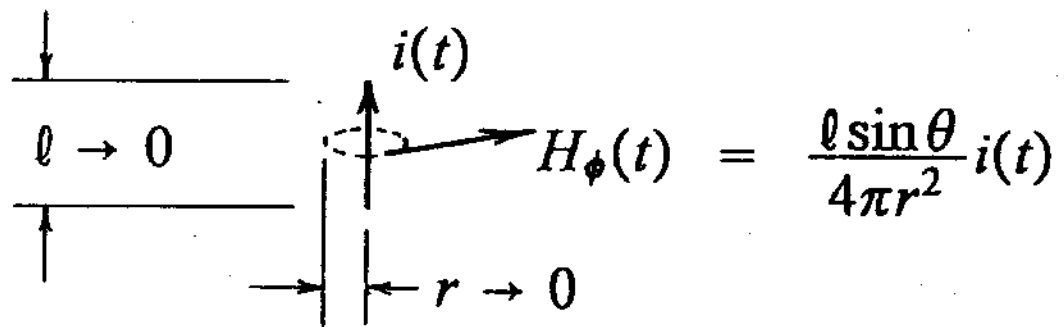
Point Current Fields

The Biot-Savart Law:



Point Current Fields

Also, for $r \rightarrow 0$, the Biot-Savart law implies



The diagram shows a horizontal current element of length l with current $i(t)$ flowing upwards. A point is located at a perpendicular distance r from the element. The magnetic field $H_\phi(t)$ is shown as a vector pointing to the right. The angle θ is between the current element and the line to the point. The distance r is the perpendicular distance from the current element to the point.

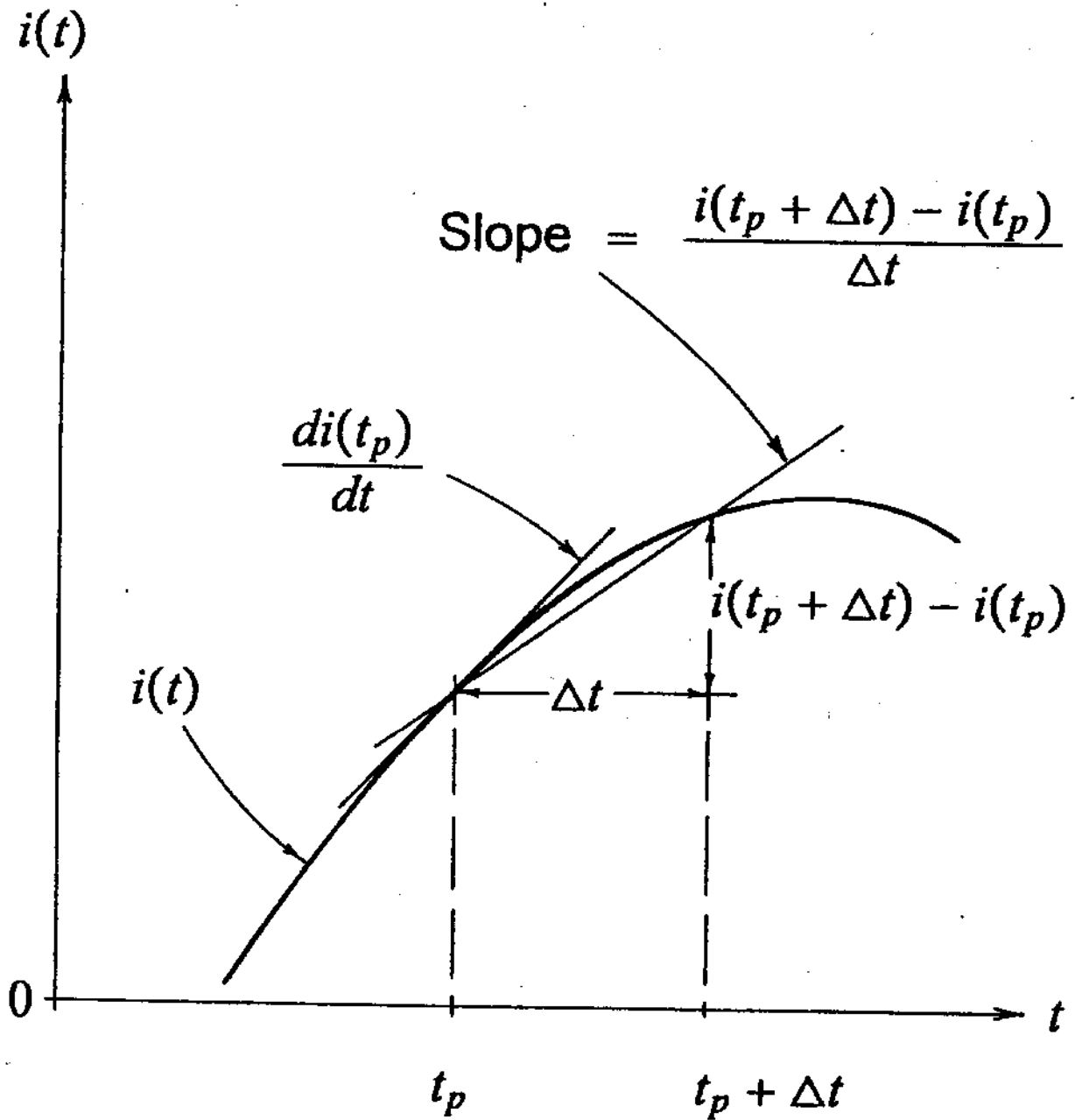
$$H_\phi(t) = \frac{l \sin \theta}{4\pi r^2} i(t)$$

And, with a propagation rate equal to c , the speed of light, it would seem that, if $r \gg 0$

$$H_\phi(t) = \frac{l \sin \theta}{4\pi r^2} i(t - r/c)$$

BUT, if a current is time-varying, $di(t)/dt \neq 0$, so charge flow increases and decreases. And that causes *radiation*. SO, when $r \gg 0$, there must be a second component of $H_\phi(t)$ that is proportional to $1/r$, and to $di(t - r/c)/dt$.

The Derivative of $i(t_p)$ and Its Definition



Point Current Fields

The derivative of $i(t)$ at t is defined to be

$$\frac{di(t)}{dt} \equiv \lim_{\Delta t \rightarrow 0} \left[\frac{i(t + \Delta t) - i(t)}{\Delta t} \right]$$

So, the derivative of $i(t)$ when $t = t - r/c$ is

$$\frac{di(t - r/c)}{dt} \equiv \lim_{\Delta t \rightarrow 0} \left[\frac{i(t - r/c + \Delta t) - i(t - r/c)}{\Delta t} \right]$$

And, replacing Δt with $r/c \rightarrow 0$, that becomes

$$\frac{di(t - r/c)}{dt} = \frac{i(t) - i(t - r/c)}{r/c}$$

which says, so long as $r/c \rightarrow 0$, or $r \rightarrow 0$,

$$i(t) = i(t - r/c) + \frac{r}{c} \frac{di(t - r/c)}{dt}$$

Point Current Fields

So, it is clear that so long as $r \rightarrow 0$,

$$\begin{aligned} H_{\phi}(t) &= \frac{\ell \sin \theta}{4\pi r^2} i(t) \\ &= \frac{\ell \sin \theta}{4\pi r^2} \left[i(t - r/c) + \frac{r}{c} \frac{di(t - r/c)}{dt} \right] \end{aligned}$$

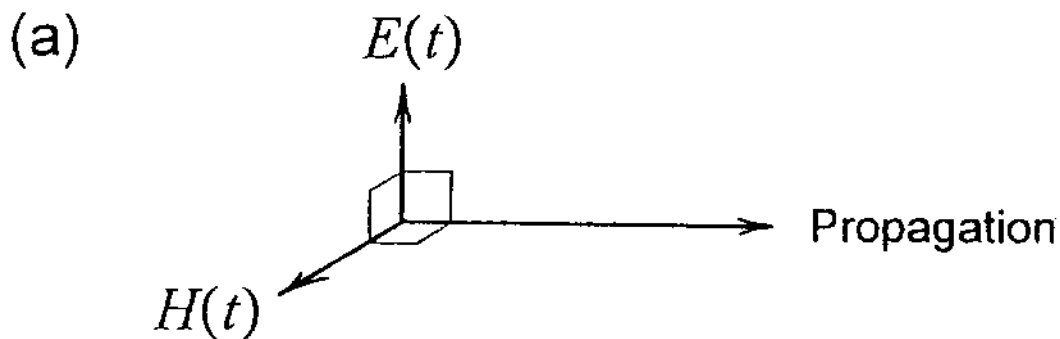
However, there is not one reason – physical, or mathematical – for the description of $H_{\phi}(t)$ to change as further propagation occurs and makes $r \gg 0$.

In other words, for all r from 0 to ∞ ,

$$H_{\phi}(t) = \frac{\ell \sin \theta}{4\pi r^2} \left[i(t - r/c) + \frac{r}{c} \frac{di(t - r/c)}{dt} \right]$$

Point Current Fields

Poynting's theorem says a radiated H-field will always be accompanied by a radiated E-field. And, those fields will be related to each other as follows:



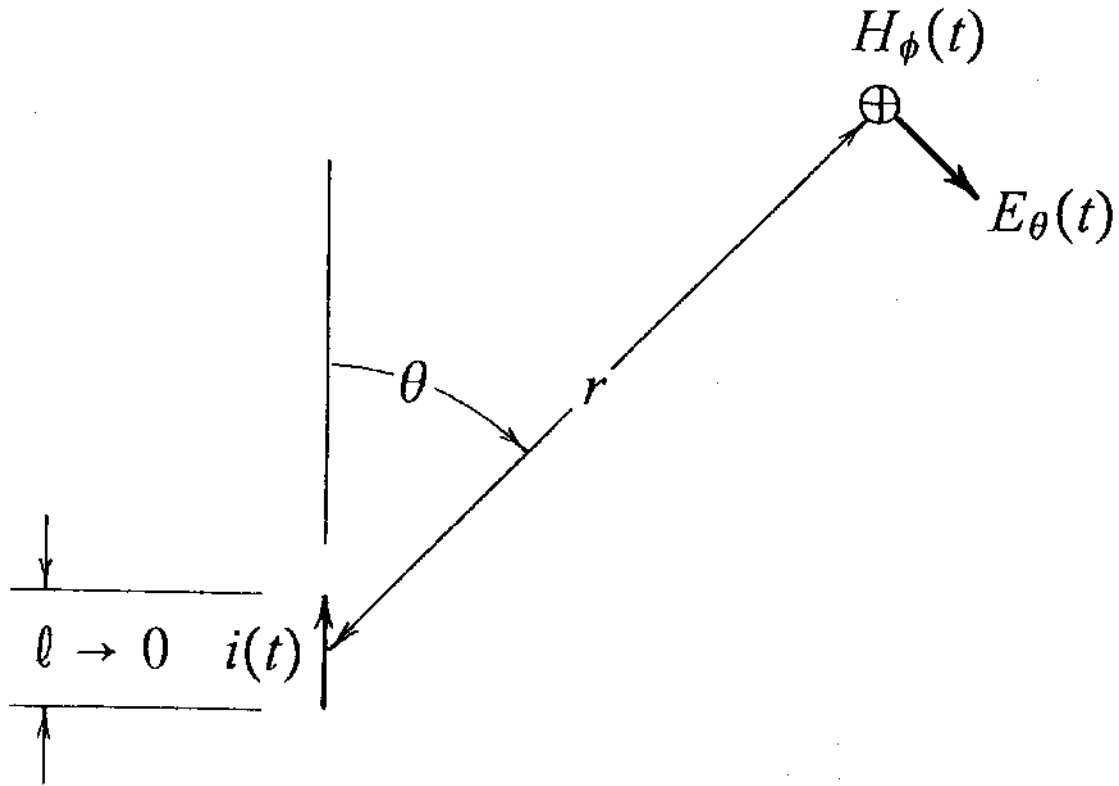
The velocity of propagation in free space is

$$c = 3 \times 10^8 \text{ meters/second}$$

(b)
$$E(t) = Z_0 H(t)$$

where $Z_0 = 120\pi$ ohms, the characteristic impedance of free space.

THE FIELDS OF A POINT CURRENT

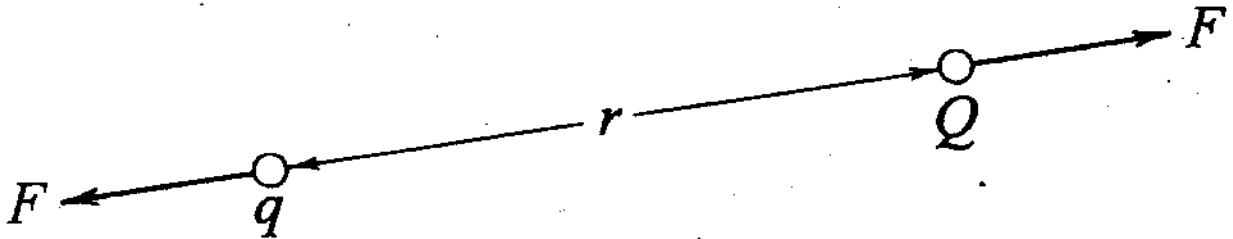


$$H_\phi(t) = \frac{\ell \sin \theta}{4\pi r^2} \left[i(t - r/c) + \frac{r}{c} \frac{\partial i(t - r/c)}{\partial t} \right]$$

$$E_\theta(t) = Z_0 \frac{\ell \sin \theta}{4\pi r^2} \left[\frac{r}{c} \frac{\partial i(t - r/c)}{\partial t} \right]$$

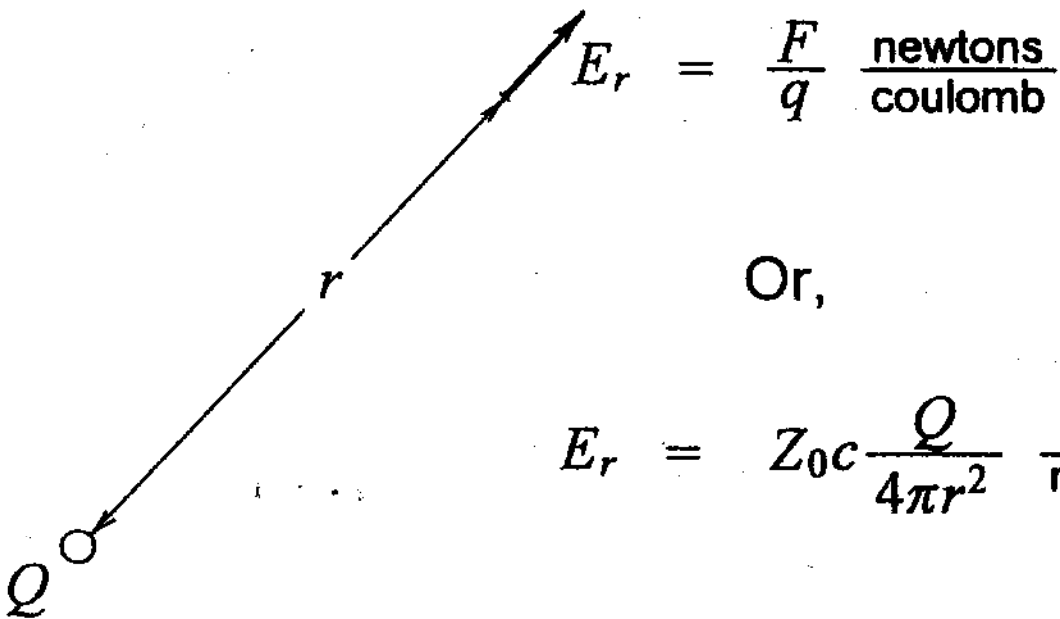
Point Charge Fields

Coulomb's Law:



$$F = \frac{1}{\epsilon_0} \frac{qQ}{4\pi r^2} = Z_0 c \frac{qQ}{4\pi r^2}$$

And, the E-field of Q is



Point Charge Fields

So, for $r \rightarrow 0$, the time-varying point charge $Q_n(t)$ has the E-field

$$E_r(t) = \frac{Z_0 c}{4\pi r^2} Q_n(t)$$

where

$$Q_n(t) = \int [i_{n-1}(t) - i_{n+1}(t)] dt$$

BUT, as with $H_\phi(t)$ and $i(t)$, if $r \gg 0$, then

$$Q_n(t) \leftarrow Q_n(t - r/c) + \frac{r}{c} \frac{\partial Q_n(t - r/c)}{\partial t}$$

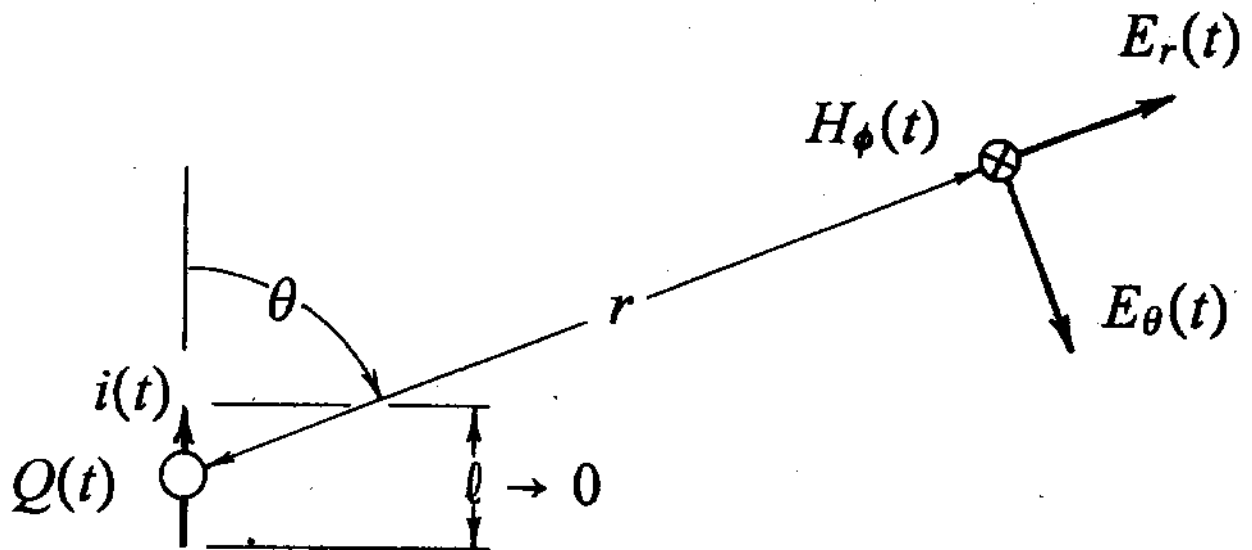
and, for all r from 0 to ∞ ,

$$E_r(t) = \frac{Z_0 c}{4\pi r^2} \left[Q_n(t - r/c) + \frac{r}{c} \frac{\partial Q_n(t - r/c)}{\partial t} \right]$$

THE CHARGE ELEMENT

$$Q(t) \circ + i(t) \uparrow \begin{array}{|c|} \hline \ell \\ \hline \end{array} \rightarrow 0 = Q(t) \circ \uparrow \begin{array}{|c|} \hline \ell \\ \hline \end{array} i(t)$$

AND ITS FIELDS



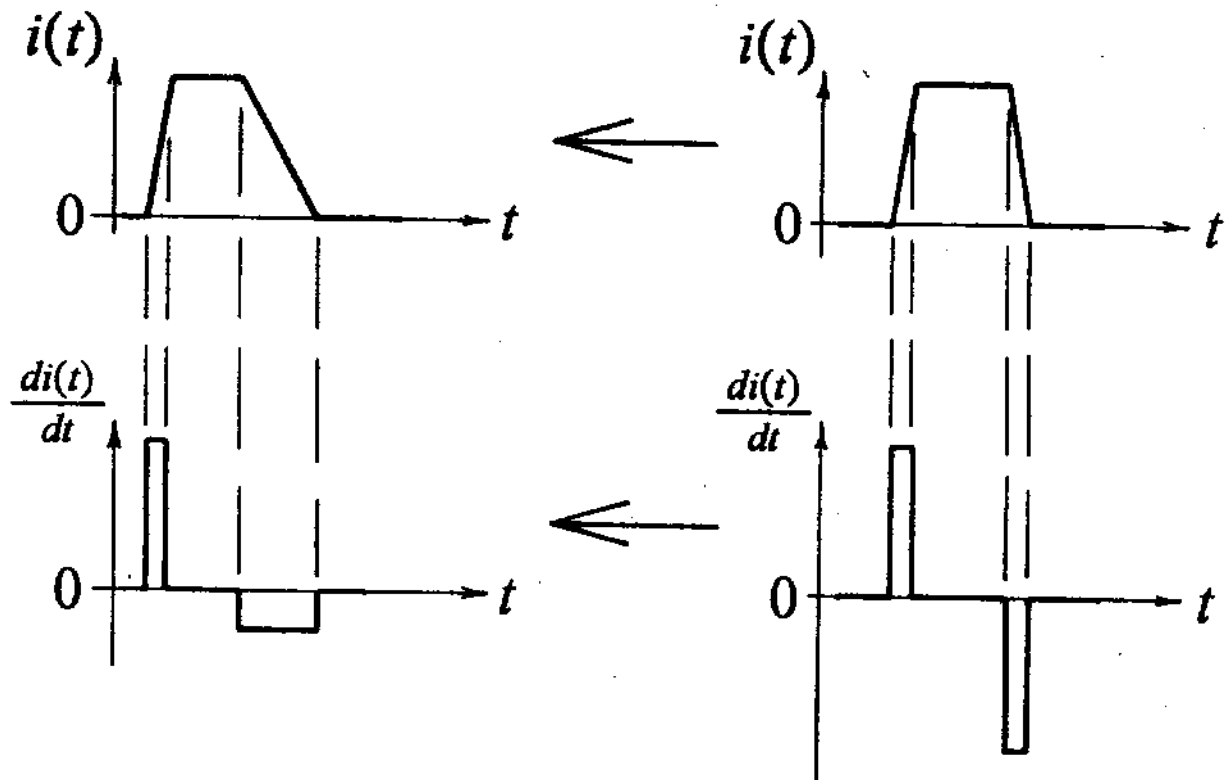
$$E_r(t) = \frac{Z_0 c}{4\pi r^2} \left[Q(t - r/c) + \frac{r}{c} \frac{\partial Q(t - r/c)}{\partial t} \right]$$

$$H_\phi(t) = \frac{\ell \sin \theta}{4\pi r^2} \left[i(t - r/c) + \frac{r}{c} \frac{\partial i(t - r/c)}{\partial t} \right]$$

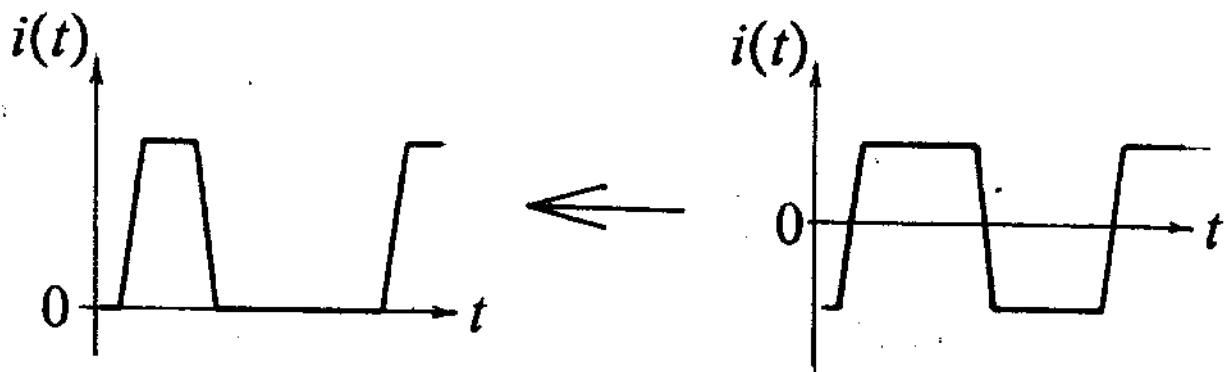
$$E_\theta(t) = Z_0 \frac{\ell \sin \theta}{4\pi r^2} \left[\frac{r}{c} \frac{\partial i(t - r/c)}{\partial t} \right]$$

Example: Minimizing Clock-Pulse EMI

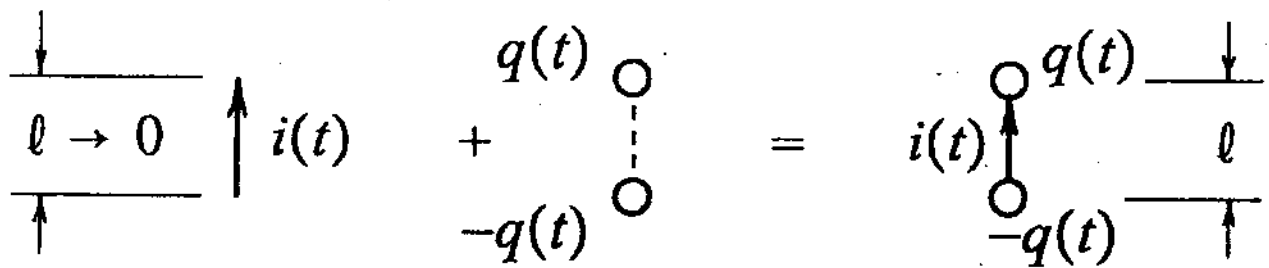
- (1) Equalize the clock current rise and fall times to minimize radiated fields.



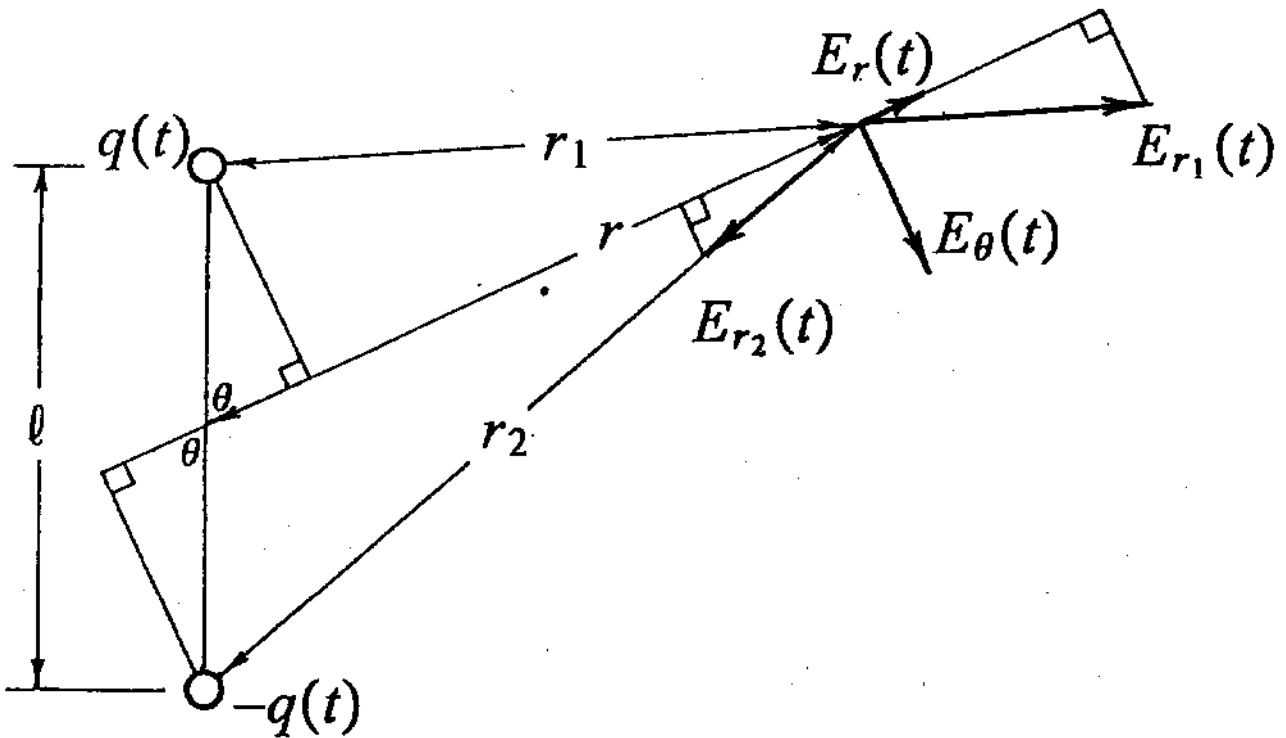
- (2) Symmetrize the current to minimize the inductive H-field.



THE CURRENT ELEMENT



. . . AND IT'S FIELDS



Since $l \rightarrow 0$ and $r \gg l$, $r_1 \cong r - \frac{l}{2} \cos \theta$,

$r_2 \cong r + \frac{l}{2} \cos \theta$, $1/r_1^n \cong 1/r^n$, $1/r_2^n \cong 1/r^n$,

and $r_2^2 - r_1^2 = (r_2 - r_1)(r_2 + r_1) \cong 2rl \cos \theta$.

Therefore, $E_r(t) \cong E_{r_1}(t) + E_{r_2}(t)$

$$\begin{aligned} &= \frac{Z_0 c}{4\pi} \left[\frac{q(t)}{r_1^2} + \frac{-q(t)}{r_2^2} \right] \\ &= \frac{Z_0 c}{4\pi} \left[\frac{r_2^2 - r_1^2}{r_1^2 r_2^2} \right] q(t) \\ &\cong Z_0 c \frac{2r\ell \cos\theta}{4\pi r^4} q(t) \end{aligned}$$

So,

$$E_r(t) = Z_0 c \frac{\ell \cos\theta}{2\pi r^3} \left[q(t - r/c) + \frac{r}{c} \frac{\partial}{\partial t} q(t - r/c) \right]$$

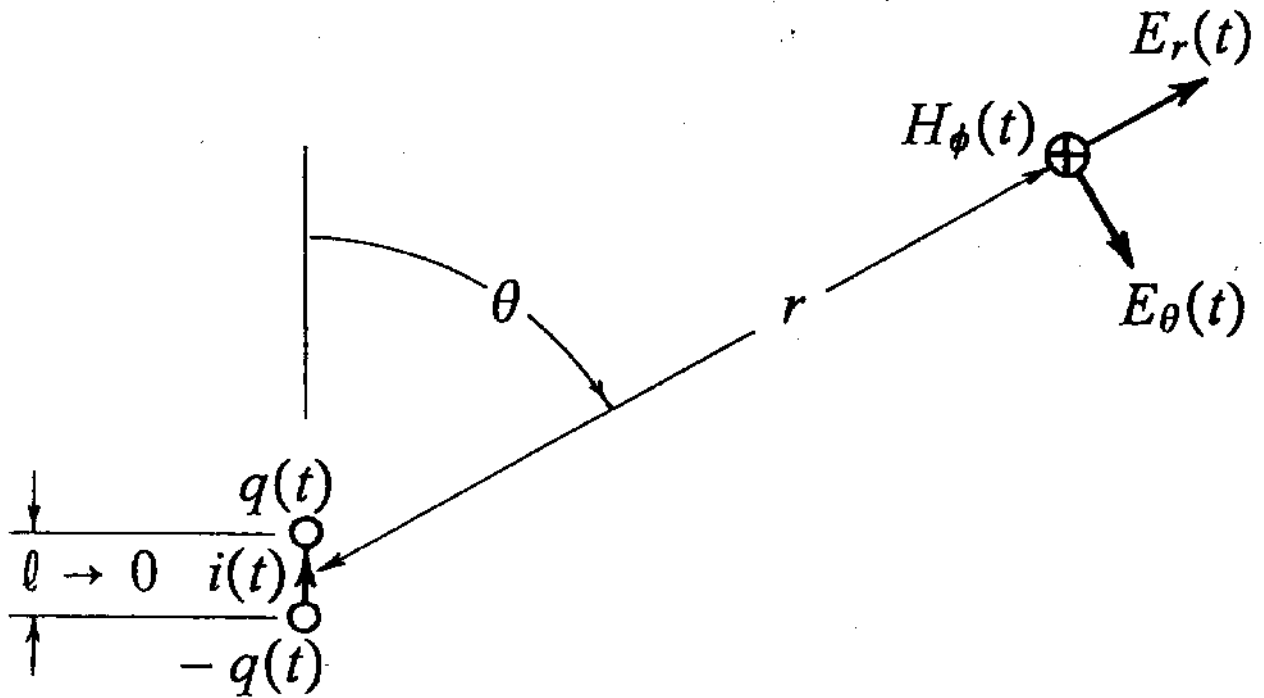
And,

$$\begin{aligned} E_\theta(t) &= \frac{\ell \sin\theta}{2r_1} E_{r_1}(t) + \frac{\ell \sin\theta}{2r_2} E_{r_2}(t) \\ &= Z_0 c \frac{\ell \sin\theta}{2} \left[\frac{q(t)}{4\pi r_1^3} + \frac{q(t)}{4\pi r_2^3} \right] \end{aligned}$$

So,

$$E_\theta(t) = Z_0 c \frac{\ell \sin\theta}{4\pi r^3} \left[q(t - r/c) + \frac{r}{c} \frac{\partial}{\partial t} q(t - r/c) \right]$$

THE CURRENT ELEMENT'S FIELDS



$$E_r(t) = Z_0 \frac{2l \cos \theta}{4\pi r^2} \left[\frac{c}{r} q(t - r/c) + \frac{\partial}{\partial t} q(t - r/c) \right]$$

$$H_\phi(t) = \frac{l \sin \theta}{4\pi r^2} \left[i(t - r/c) + \frac{r}{c} \frac{\partial}{\partial t} i(t - r/c) \right]$$

$$E_\theta(t) = Z_0 \frac{l \sin \theta}{4\pi r^2} \left[\frac{c}{r} q(t - r/c) + \frac{\partial}{\partial t} q(t - r/c) \right]$$

$$+ Z_0 \frac{l \sin \theta}{4\pi r^2} \left[\frac{r}{c} \frac{\partial}{\partial t} i(t - r/c) \right]$$

EQUIVALENCE OF CHARGE ELEMENTS AND CURRENT ELEMENTS

